

Quantum equation of tachyon

D.V. Filippov

RECOM Russian research Center “Kurchatov Institute”, **Russia**

www.uf.narod.ru

Quantum equation of tachyon

As discussed in detail by **Recami&Mignani** [R. Mignani, E. Recami // *Nuovo Cimento*, 30A, #4, 533 (1975); *Lett. Nuovo Cimento*, 11, #8, 417 (1974).], the existence of the **tachyon** (i.e. a superluminal particle) has no conflict with the theory of special relativity.

Within the framework of **classical electrodynamics**, **Recami** has shown that **superluminal electric charges** (tachyons) can be described similarly to **subluminal magnetically charged** particles.

Meanwhile, **G. Lochak** paid attention to the fact that the Tushek-Salam gauge invariance for the **massless Dirac equation** gives rise to an equation describing the **magnetic monopoles** [G.

Lochak // *Ann. Fond. L.de Broglie*, 1983, #8, 345; *Ann. Fond. L.de Broglie*, 1984, #9, 5.]



Quantum equation of tachyon

The new equation which is proposed

1. obeys the local Tushek-Salam gauge invariance with allowance for the **linear mass term**, that is, it represents a development of **Lochak's magnetic monopole** theory, and
2. describes the magnetic monopoles that coincide, in the classical approximation, with the **Recami tachyons**.



Quantum equation of tachyon

$$\left\{ \gamma_{\mu} \left(\partial_{\mu} - \frac{\mathbf{g}}{\hbar \mathbf{c}} \gamma_5 \cdot \mathbf{B}_{\mu} \right) + \frac{m\mathbf{c}}{\hbar} \Gamma \right\} \psi = \mathbf{0},$$

$$\bar{\psi} \left\{ \gamma_{\mu} \left(-\overleftarrow{\partial}_{\mu} - \frac{\mathbf{g}}{\hbar \mathbf{c}} \gamma_5 \cdot \mathbf{B}_{\mu} \right) + \bar{\psi} \cdot \frac{m\mathbf{c}}{\hbar} \Gamma \right\} = \mathbf{0},$$

$$\gamma_5 \Gamma = -\Gamma \gamma_5,$$

$$\Gamma = \mathbf{a}_k \gamma_k \cdot \gamma_5, \quad \mathbf{a}_k \mathbf{a}_k = \mathbf{1},$$

$$\Gamma^2 = -\mathbf{a}_k \gamma_k \cdot \mathbf{a}_m \gamma_m = -\mathbf{1},$$



Quantum equation of tachyon

Each of the equations splits into two independent equations in the **Weyl** representation:

$$\left[\frac{1}{c} \frac{\partial}{\partial t} - \vec{\sigma} \cdot \nabla - i \frac{g}{\hbar c} (\chi + \vec{\sigma} \vec{B}) - i \frac{mc}{\hbar} \right] \xi = 0$$

$$\left[\frac{1}{c} \frac{\partial}{\partial t} + \vec{\sigma} \cdot \nabla + i \frac{g}{\hbar c} (\chi - \vec{\sigma} \vec{B}) + i \frac{mc}{\hbar} \right] \eta = 0$$



Quantum equation of tachyon

The fact that this equation really describes the tachyons becomes evident from the **dispersion characteristic** equation:

$$\left(\frac{\omega}{c} - \frac{mc}{\hbar} \right)^2 = \vec{k}^2$$

Dirac:
$$\left(\frac{\omega}{c} \right)^2 = k_0^2 + \vec{k}^2$$



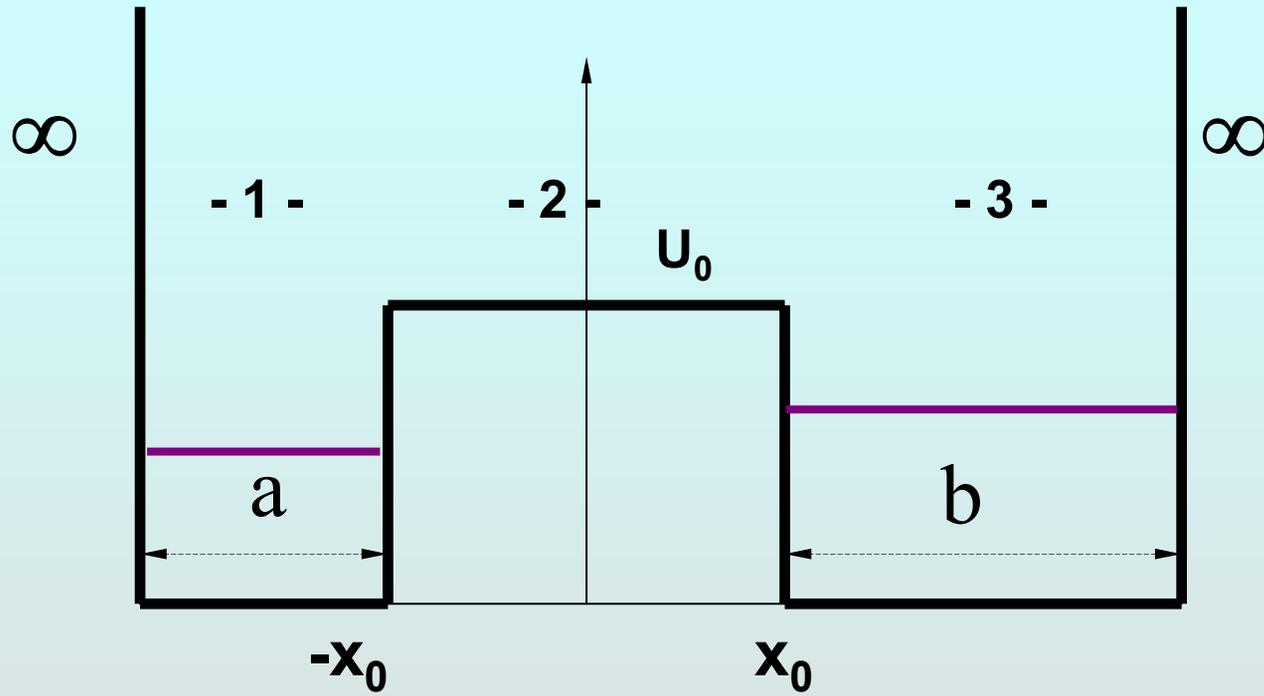
The effect of quantum tunnel resonance

D.V. Filippov

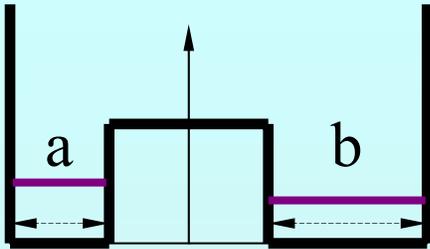
RECOM Russian research Center “Kurchatov Institute”, **Russia**

www.uf.narod.ru

The effect of quantum tunnel resonance



The effect of quantum tunnel resonance



$$i\hbar \frac{\partial}{\partial t} \tilde{\psi}(t, \mathbf{x}) = \hat{H} \tilde{\psi} = \left[-\frac{\hbar^2}{2m} \cdot \frac{\partial^2}{\partial \mathbf{x}^2} + U(\mathbf{x}) \right] \tilde{\psi}$$

$$\tilde{\psi}(t, \mathbf{x}) = \exp(-iEt/\hbar) \cdot \psi(\mathbf{x})$$

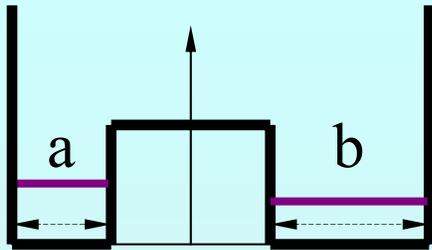
$$\psi_1 = \psi_0 \cdot \sin[k(\mathbf{x} + \mathbf{x}_0 + \mathbf{a})]$$

$$\psi_2 = \psi_0 \cdot [B_- \exp(-\kappa \mathbf{x}) - B_+ \exp(\kappa \cdot \mathbf{x})]$$

$$\psi_3 = A \cdot \psi_0 \cdot \sin[k(\mathbf{x} - \mathbf{x}_0 - \mathbf{b})]$$



The effect of quantum tunnel resonance



$$f \equiv -\frac{k}{\kappa} \operatorname{ctg}(ka) \quad g \equiv -\frac{k}{\kappa} \operatorname{ctg}(kb)$$

$$\kappa X_0 \gg 1$$

$$A = \frac{\pm \sqrt{k^2 + (\kappa \cdot g)^2}}{\sqrt{k^2 + (\kappa \cdot f)^2}} \cdot [\operatorname{ch}(2 \cdot \kappa X_0) - f \cdot \operatorname{sh}(2 \cdot \kappa X_0)]$$

In resonance case: $ka=kb + \pi \cdot n$, i.e. $g=f$.

We received *exact* decision $A=\pm 1$, that is the particle equiprobably occupies both "holes".

We shall notice, that the condition of a resonance depends only on a ratio of "width" of left and right "hole" and

does not depend on width and height of "barrier"

between them. The allowable width of a resonance depends on the sizes of "barrier".



Study of the gas released upon electric explosion of titanium foils in liquids

**Urutskoev L.I.¹, Govorun A.P.¹, Gulyaev A.A.¹,
Demkin S.A.², Dorovskoi V.M.², Elesin L.A.²,
Kuznetsov V.L.², Petrushko S.V.¹, Steblevskii A.V.³,
Stolyarov V.L.², Fedotov V.G.⁴, Filippov D.V.¹**

1. RECOM Russian research Center “Kurchatov Institute”, Russia

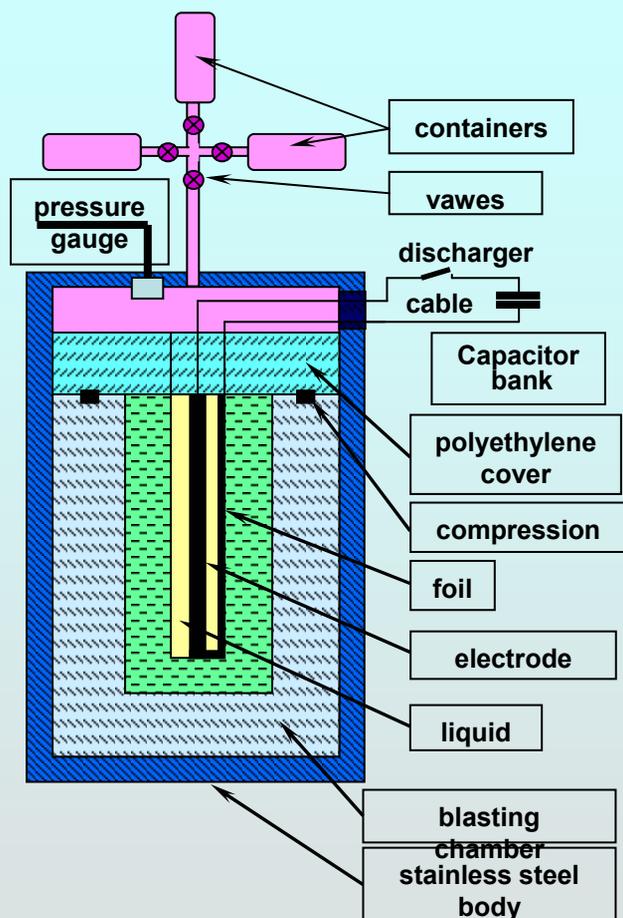
2. Kurchatov Institute

3. Institute of Inorganic Chemistry

4. Institute of Chemical Physics

www.uf.narod.ru

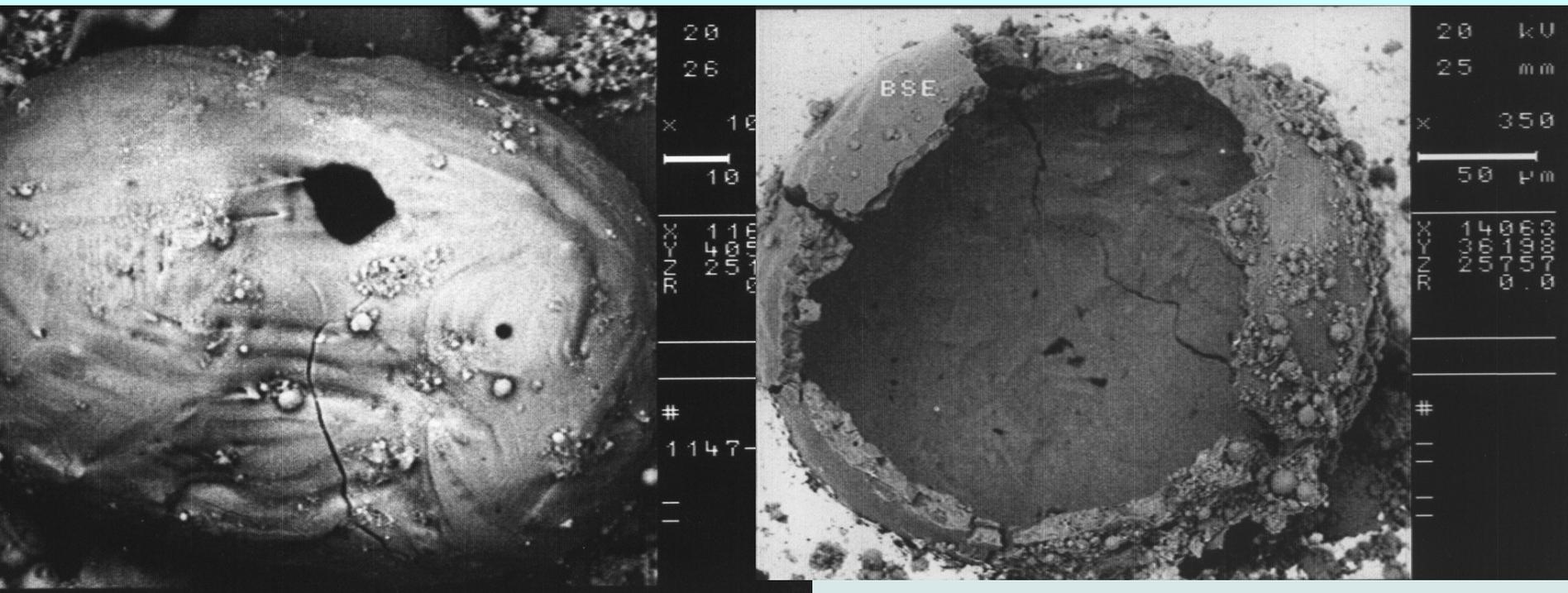
Study of the gas released upon electric explosion of titanium foils in liquids



1. After an electric explosion of titanium foil in water, the sediment formed contains spherical hollow particles filled with H_2 . In some particles, hydrogen pressure is so high that the particles are destroyed.
2. Noteworthy is the absence of O_2 in hollow particles.
3. An inhomogeneous content of admixtures in Ti has been detected over different topological fragments of the sample.



Study of the gas released upon electric explosion of titanium foils in liquids



1. After an electric explosion of titanium foil in water, the sediment formed contains spherical hollow particles filled with H_2 . In some particles, hydrogen pressure is so high that the particles are destroyed.
2. Noteworthy is the absence of O_2 in hollow particles.
3. An inhomogeneous content of admixtures in Ti has been detected over different topological fragments of the sample.